

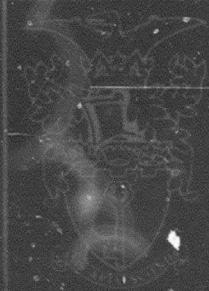
18233

TR 69275

DECEMBER

1969

AD705024



Crown Copyright
1969

ROYAL AIRCRAFT ESTABLISHMENT

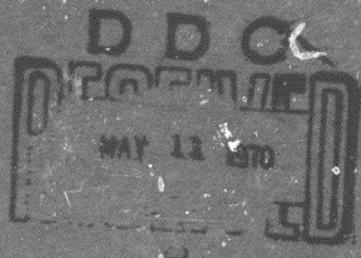
TECHNICAL REPORT 69275

THE ORBIT OF ARIEL 3

(1967-42A)

by

R. H. Gooding



Reproduced by the
CLEARINGHOUSE
for Federal Scientific & Technical
Information Springfield Va. 22151

UNLIMITED

MINISTRY OF TECHNOLOGY
FARNBOROUGH HANTS

U.D.C. 521.6 : 629.19.077.3 : 629.195

ROYAL AIRCRAFT ESTABLISHMENT

Technical Report 69275

December 1969

THE ORBIT OF ARIEL 3 (1967-42A)

by

R. H. Gooding

SUMMARY

The definitive orbit for Ariel 3 has been computed, from Minitrack observations, for a period of 27½ months from the launch of the satellite. The orbit was represented by a model with seven independent orbital parameters and the values of these parameters were determined, and are listed, at three-day intervals. Typical accuracies are 10^{-5} in eccentricity and 4" in angular parameters, that is, about $\frac{1}{2}$ km in position.

A curious feature of the secular variation of orbital inclination, viz. that the expected decrease of about 0.02° appeared to occur over a three-month period instead of the full 27½-month period, is discussed but has not been explained.

Departmental Reference: Space 327

CONTENTS

	<u>Page</u>
1 INTRODUCTION	3
2 OBSERVATIONS	3
3 ORBITAL MODEL	4
4 RESULTS	6
4.1 Main results	6
4.2 Results involving the parameter M_3	9
4.3 Results involving Hewitt camera observations	10
5 ACCURACY OF POSITION COMPUTATION	10
6 DISCUSSION	12
7 CONCLUSIONS	16
Acknowledgement	16
Table 1 Minitrack stations observing Ariel 3	17
Table 2 Orbital parameters of Ariel 3	19
References	26
Illustrations	Figures 1-3
Detachable abstract cards	-

1 INTRODUCTION

Ariel 3, the third of the series of satellites being launched in the scientific programme of Anglo-American co-operation, was the first spacecraft to be built entirely in Britain. The five experiments in the payload were concerned¹ with electron density and temperature (University of Birmingham), VLF radio waves (University of Sheffield), cosmic radio noise (Nuffield Radio Astrohometry Laboratory, Jodrell Bank), molecular oxygen (Meteorological Office) and terrestrial noise sources (Radio and Space Research Station).

Known as S53 or UK3 before launch, and as Ariel 3 or 1967-42A afterwards, the satellite was placed in a near-polar, near-circular orbit at 18^h UT on 5 May 1967 by a Scout rocket launched from the Western Test Range, California.

The definitive orbit of Ariel 3, as for Ariel 2², has been derived at R.A.E. from Minitrack (interferometer) data provided by NASA. Orbital parameters have been obtained at three-day intervals by the use of the new computer program PROP³. They are tabulated, for the first 27½ months of the satellite's lifetime, in this Report. For epochs up to 1968 JAN 16 the work was done on London University's Atlas computer and for subsequent epochs it was done on the ICL 1907 computer at R.A.E.

2 OBSERVATIONS

The STADAN Minitrack network now consists of ten stations - five of the original twelve stations² have disappeared and there are three new ones. These are listed in Table 1, with their assumed positions in standard geocentric co-ordinates (x-axis towards the Greenwich meridian).

Observations consist of pairs of direction cosines. Their a priori accuracy (s.d.) has been taken, as usual^{3,4}, at 0.00029, equivalent to 1' in angular measure, though their true accuracy is believed to be worse than this.

Times are given in the UTC system, i.e. the system defined by WWV time transmissions from America, and have not been corrected during the orbital determination, except that the times of observations made on 1 February 1968 and used at epoch 1968 JAN 31 had to be reduced by 0.1^s to allow for the step advance of UTC at Feb 1.0. Times should be accurate to about 1 ms; no allowance for timing error was made by PROP.

In total, about 10000 Minitrack observations were used, covering the period from launch until about a week before the satellite's transmitter was switched off (1 September 1969); i.e. these were about 12 per day. They were received from NASA on punched cards, suitable for direct input to PROP. The epochs for orbit determination were taken at three-day intervals, and always at midnights (unlike the Ariel 2 epochs which were at ascending nodes). The orbit determination at each epoch used observations over a four-day period, allowing one-day overlaps in the periods of validity of the resulting orbital parameters, but observations were not (in general) used twice; observations on an 'overlap day' were divided into two sets, by alternate allocation, for use with the epochs before and after the overlap day. About 650 of the observations were rejected during analysis, but this includes (a) a rather high rejection rate during the first 3½ months - all the observations from Orroral were being rejected at one stage - and (b) nearly 100 observations, over a period of a month (Jan-Feb, 1968), which all had a one wavelength error in the north-south direction cosine due to a temporary error in the NASA program⁵ for processing the raw data; the normal rejection rate was about 4½%.

The number of Minitrack observations per day varied, of course, but there was at least one on every day of the period covered, apart from the week 28 November-4 December 1967, for which there was a complete absence of data.

A few observations of Ariel 3 were made by the Hewitt camera⁶ at Malvern. Among these, 8 observations came from a pass on the evening of 10 April 1969 and 12 observations came from two passes close to midnight on 19 June 1969 and 21 June 1969. It was decided to incorporate these 20 observations into re-runs of the orbit determinations at the appropriate two epochs, to see how fit and accuracy were affected. The remaining Hewitt camera observations, and the many visual observations of Ariel 3, have not been used.

3 ORBITAL MODEL

The orbital model of the program PROP is not the same as that used in the analysis of the orbit of Ariel 2. Eccentricity, inclinations etc. are defined slightly differently in the two programs, and the connecting relations are given in Appendix C of Ref.3.

The model allows some choice as to the set of orbital parameters which represent the orbit and which are determined from fitting to observations. The set chosen for the Ariel 3 orbit contained seven parameters, viz., e_0 (eccentricity), i_0 (inclination), Ω_0 (right ascension of the node), ω_0 (argument of perigee), M_0 (mean anomaly), M_1 (mean motion) and M_2 (half the mean acceleration). The first four parameters are epoch values of mean elements, as defined in Refs. 7 and 8 and the last three are the coefficients in the polynomial representation of (mean) mean anomaly:

$$M = M_0 + M_1 t + M_2 t^2 ,$$

where t is measured from epoch.

Secular rates of change of e , i , Ω and ω (i.e. the polynomial coefficients e_1 , i_1 , Ω_1 and ω_1) were computed, inside PROP at the beginning of each iteration of the differential-correction process, as functions of the seven independent parameters^{3,8}. These quantities, together with the long-periodic and short-periodic terms computed at each observation time, represented orbit perturbations due to drag and to the earth's zonal harmonics up to J_9 . The along-track effect of the tesseral harmonic $J_{2,2}$ was represented as usual^{3,8} using the value 1.8×10^{-6} , but, apart from this, tesseral harmonics were neglected. Luni-solar perturbations were ignored; their effect on Ariel 3 over a period of two or three days from any epoch is very small.

The decision not to have an eighth parameter M_3 , which would have made the M polynomial a cubic, was justified by some test runs, early in the lifetime, which showed that no significant improvement in fit would result and that the value of M_3 itself would not be significant. For the two epochs during the week of missing data, however, M_3 was included in the model and a reasonable fit thereby obtained to data before and after the gap, covering a period of $10\frac{1}{2}$ days. In retrospect, the decision to omit M_3 is open to question because drag increased fairly steadily through the $27\frac{1}{2}$ months considered (as the values of M_2 in Table 2 show), the maximum effects being at the end of March 1969; repetition of two of the runs, with M_3 included, showed that at this stage significant improvement in fit would be obtained (see section 4).

Apart from the omission of M_3 during periods of high drag, the main limitation of the orbital model is in the neglect of important tesseral-harmonic perturbations and, in particular, of the perturbation in inclination due to $J_{2,2}$. This perturbation has a period of just under 12 hours and an amplitude of about $0^{\circ}.002$, equivalent to a maximum position error of about $\frac{1}{2}$ km.

4 RESULTS

4.1 Main results

The orbital parameters obtained from the orbit determinations, together with certain additional information, are listed in Table 2. Successive columns of the table provide the following quantities, zero suffixes being omitted from a_0 etc:-

Epoch date (0^h UTC understood).

Semi-major axis, a (km).

Eccentricity, e .

Perigee height, h_p (km).

Inclination, i (degrees).

Right ascension of the node, Ω (degrees).

Argument of perigee, ω (degrees).

Mean argument of latitudue, $M_0 + \omega$ (degrees).

Mean motion, M_1 (degrees/day).

Half acceleration, M_2 (degrees/day²).

Number of observations used, N .

Number of observations rejected, K .

Extent of the observations, D (days).

Standard deviation of an observation of unit weight, ϵ .

Modified Julian Day number of epoch date, MJD.

The orbital parameters are the seven quantities $e, i, \Omega, \omega, M_0, M_1$ and M_2 , but $M_0 + \omega$ is given instead of M_0 because of the high correlation between M_0 and ω . This correlation arises directly from the fact that the orbit is so nearly circular, and Ref.3 may be consulted for further explanation. (The appropriate value of the control parameter JELTYP³ was used to give the variance of $M_0 + \omega$ directly.)

The semi-major axis, a , is the mean element, as used by Merson⁷, defined from M_1 by

$$a = (\mu/M_1)^{\frac{1}{3}} - \frac{1}{4} J_2 R^2 (\mu/M_1)^{-\frac{1}{3}} (2 - 3 \sin^2 i) (1 - e^2)^{-\frac{1}{2}},$$

where μ is the earth's gravitational constant, J_2 is its second zonal harmonic coefficient and R is its mean equatorial radius, the values from Ref.3 being used.

The perigee height is given by

$$h_p = r_p - R_p,$$

where*

$$r_p = a(1 - e) + \frac{J_2 R^2}{4a(1 - e^2)} \left\{ \sin^2 i \cos 2\omega - (2 - 3 \sin^2 i) \left(1 + \frac{e}{1 + \sqrt{1 - e^2}} - \frac{\sqrt{1 - e^2}}{1 + e} \right) \right\}$$

and

$$R_p = R - 21.379 \sin^2 i \sin^2 \omega.$$

The right ascension of the node is nominally referred to the standard PROP equinox³, but contains a small error due to the fact that the times are given in UTC and no correction to UT1 has been made. To correct Ω to the time PROP equinox (epoch date still understood to be 0^h UTC) add $0^h.004 \times (\text{UT1} - \text{UTC})$, where the time difference is in seconds.

The 'number of observations used' includes the number rejected; i.e. the parameters have been determined, in the end, from $N - K$ observations.

After nine of the tabulated quantities - the seven orbital parameters plus semi-major axis and perigee height - are given their computed standard

*The difference between r_p and $a(1 - e)$ is important. Thus, though r_p (or rather h_p) is the right parameter to use when correlating drag behaviour with perigee height, $a(1 - e)$ is the right parameter to work with when studying the effects of the earth's odd harmonics. For Ariel 3 the difference is approximately $1.54 \cos 2\omega$ km, and it was the use of r_p instead of $a(1 - e)$ which led to the apparent discrepancy mentioned in section 3.20 of Ref.13.

deviations, to one or two significant figures, the unit in each case being that of the final figure quoted for the main quantity. Every standard deviation includes ϵ as a factor, where ϵ is given by

$$\epsilon = \{\sum (\text{Res}/0.00029)^2/(2N - 2K - 7)\}^{1/2};$$

here the summation is over all residuals, Res, in the N - K accepted east-west and north-south direction cosines, and 0.00029 is the a priori accuracy referred to in section 2. Since the actual accuracy is worse than this a priori figure⁴, the values of ϵ in Table 2 are expected to be - and are - larger than unity.

Table 2 was obtained as direct computer output from a program known as TOP (Tabulation of Orbital Parameters). This program takes, as part of its input data, the punched card output from PROP runs, so there should not be any errors in the table.

The secular rates of change e_1 , i_1 , Ω_1 and ω_1 are not given in Table 2, since they are computed internally by PROP as part of the model. It is remarked, however, that the computation of the J_2^2 component of ω_1 contained an error until the PROP3 version of the program was introduced at the end of January, 1969. PROP2, which had this error in ω_1 , was able to compensate for it almost exactly, by fitting a slightly wrong value of M_1 , and this was one reason why the error was for a long time undiscovered. To correct the results from the PROP2 runs it was only necessary to correct M_1 by an amount equal to the error in ω_1 , and a special program was written to do this. The only reason for mentioning this point is that the values of M_1 in Table 2 are the corrected values, and so are different from the values provided in the first four provisional lists of Ariel 3 parameters to be issued. (To avoid having to ask AWRE Foulness to make a correction to the Ariel 3 telemetry data analysis program after PROP3 had been introduced, it was decided to continue sending incorrect values to AWRE, by adding the appropriate deliberate error to M_1 .)

Fig.1 gives a plot of orbital inclination, each value being represented by a vertical line, two standard deviations in length, centred on the fitted value. Fig.2 gives a plot of eccentricity, but most of the time the scale is too small for standard deviations to be shown. Fig.3 shows a short section of the eccentricity curve (covering just over half a period of the perigee) with

the scale expanded sufficiently for the standard deviations to be indicated as on the inclination plot. Figs. 1 and 2 give, in addition to the definitive inclinations and eccentricities obtained at R.A.E., the SDC (NORAD) values published in Spacetrack bulletins.

4.2 Results involving the parameter M_3

For two of the runs covered by Table 2 the parameter M_3 was included in the orbital model, namely, for those of epochs 1967 NOV 29 and 1967 DEC 2. These were the epochs which occurred during the week when no Minitrack data were supplied. Without M_3 the fit was twice as bad (ϵ 5.7 for the first epoch instead of 2.8), due to the number of days spanned by the observations. The value of M_3 , omitted from Table 2 to retain a regular format, was -0.00069, with standard deviation 0.00003, for both epochs; the same observations were used in both runs, so the second set of elements is really just the first set advanced three days, with small variations in the residuals (and hence a small change in ϵ) due to limitations of the orbital model.

It was stated in section 3 that, for epochs early in the satellite's lifetime, general introduction of M_3 would not have helped. To illustrate, the complete set of parameters, when M_3 is included, for 1967 DEC 20 is as follows (with last-figure standard deviations in brackets):-

$e = 0.007329$ (15), $i = 80.1802$ (19), $\Omega = 239.0276$ (19), $\omega = 155.22$ (10),
 $M_0 + \omega = 91.7171$ (15), $M_1 = 5433.0021$ (22), $M_2 = 0.0701$ (9) and
 $M_3 = 0.0007$ (9); the value of ϵ , viz. 3.6, was actually larger (unrounded value, 3.553 as against 3.546) than for the run without M_3 , due to the loss of a degree of freedom. For certain epochs later in the lifetime, however, introduction of M_3 would have led to better fits. This may be illustrated by considering the two worst fits obtained, namely, for epochs 1969 MAR 20 (ϵ of 5.0 in Table 2) and 1969 MAR 23 (ϵ of 5.1); on repeating these runs, with M_3 included, the following results were obtained:- for 1969 MAR 20,
 $e = 0.006902$ (45), $i = 80.1665$ (20), $\Omega = 10.3404$ (19), $\omega = 124.14$ (17),
 $M_0 + \omega = 189.4465$ (46), $M_1 = 5475.7830$ (24), $M_2 = 0.1282$ (9) and
 $M_3 = 0.0077$ (9), with $\epsilon = 3.1$; for 1969 MAR 23, $e = 0.006936$ (49),
 $i = 80.1677$ (30), $\Omega = 6.4259$ (22), $\omega = 116.38$ (16), $M_0 + \omega = 48.2682$ (47),
 $M_1 = 5476.6144$ (32), $M_2 = 0.1444$ (12) and $M_3 = 0.0085$ (13), with
 $\epsilon = 3.6$.

4.3 Results involving Hewitt camera observations

The runs at epochs 1969 APR 10 and 1969 JUN 21 were repeated, with (respectively 8 and 12) Hewitt camera observations included. The following results were obtained:- for 1969 APR 10, $e = 0.006778$ (7), $i = 80.1671$ (12), $\Omega = 342.9333$ (13), $\omega = 67.28$ (12), $M_0 + \omega = 326.4188$ (10), $M_1 = 5480.4510$ (8) and $M_2 = 0.0850$ (7), with $\epsilon = 2.0$ and only the same Minitrack observation rejected as was originally rejected; for 1969 JUN 21, $e = 0.005403$ (8), $i = 80.1610$ (10), $\Omega = 248.6570$ (13), $\omega = 175.80$ (6), $M_0 + \omega = 120.8871$ (10), $M_1 = 5490.2793$ (7) and $M_2 = 0.0584$ (3), with $\epsilon = 2.4$ and two of the Minitrack observations rejected that had previously been accepted. On comparison with corresponding entries in Table 2 it may be seen that for the first run there is little change - the maximum change in a parameter is for $M_0 + \omega$, the change being about twice the original standard deviation, and no standard deviation has decreased by a factor of more than $1\frac{1}{2}$; for the other run, however, there is a large change in eccentricity, nearly five times the original standard deviation, and the standard deviations for e , ω and $M_0 + \omega$ have all been reduced by factors of more than 2. (It is worth remarking that the change in e was caused entirely by the introduction of the Hewitt camera observations, and not at all by the subsequent rejection of two Minitrack observations.)

A reasonable conclusion is that Hewitt camera observations, of high accuracy, are compatible with Minitrack observations of poorer accuracy. For a high-inclination satellite like Ariel 3 the effect is not very significant if the Hewitt camera observations all come from a single pass, but there is a great improvement in accuracy when observations from two or more passes are available.

5 ACCURACY OF POSITION COMPUTATION

As with Ariel 2 it was required, for correlation with on-board experiments, that the definitive orbital parameters should be good enough for position to be computable from them to better than 1 km. In the paper on the Ariel 2 orbit² the accuracy of position computation was considered by reference to plots of $\{\sigma^2(x) + \sigma^2(y) + \sigma^2(z)\}^{\frac{1}{2}}$, where the variances $\sigma^2(x)$, $\sigma^2(y)$ and $\sigma^2(z)$ are functions of time and the covariance matrix of the orbital parameters, and by comparison of such plots with plots, during overlap periods, of $\{(x_1 - x_2)^2 + (y_1 - y_2)^2 + (z_1 - z_2)^2\}^{\frac{1}{2}}$, where x_1, y_1, z_1 denote satellite

co-ordinates computed from orbital parameters at the epoch before the given overlap period and x_2, y_2, z_2 denote co-ordinates computed from parameters at the following epoch. This approach would have been equally possible for Ariel 3, using the program PREP⁸, but it was decided that it would be adequate to consider the question by looking directly at standard deviations of orbital parameters and interpreting these as maximum position errors after 1½ days.

The justification for this approach is that, with $M_o + \omega$ rather than M_o taken as a parameter, large correlations between parameters did not occur. (Occasional correlations as large as ± 0.4 occurred, usually involving e and one other parameter.) For a satellite, like Ariel 3, in an orbit which is nearly circular and not too far from polar, the maximum effects, on position after 1½ days, of one-sigma errors in the parameters are approximately as follows: $2 \sigma(e)$, $\sigma(i)$, $\sigma(\Omega)$, $2 \sigma(e) \sigma(\omega)$, $\sigma(M_o + \omega)$, $1\frac{1}{2} \sigma(M_1)$ and $2 \sigma(M_2)$, where angle sigmas are now taken to be in radians. The main effects here of e , ω and $M_o + \omega$ are the along-track errors which arise from the expression of argument of latitude in the form

$$u = (M + \omega) + 2e \sin \{(M + \omega) - \omega\} .$$

Let us consider 'maximum position effects' for two different sets of sigmas: first, the maximum value of each sigma that occurs anywhere in Table 2, and, second, maximum values during, roughly speaking, the best ninety per cent of the time. Denoting sigmas from the two sets by σ_1 and σ_2 respectively (with angles in degrees again), and the corresponding maximum position effects, in km, by MPE_1 and MPE_2 , we have the following table:-

	σ_1	MPE_1	σ_2	MPE_2
e	0.000072	1.1	0.000020	0.3
i	0.0043	0.5	0.0030	0.4
Ω	0.0048	0.6	0.0025	0.3
ω	0.51	0.8	0.17	0.3
$M_o + \omega$	0.0074	0.9	0.0031	0.4
M_1	0.0043	0.8	0.0023	0.4
M_2	0.0019	0.5	0.0019	0.5

Since $(\sum MPE_2^2)^{\frac{1}{2}} = 1.0$ km it is reasonable to claim that the accuracy requirements are met most of the time. If we consider the accuracy of height only, i.e. of $r = a \{1 - e \cos (M + \omega - \omega_0)\}$, then only the maximum one-sigma effects $a \sigma(e)$ and $a e \sigma(\omega)$ are significant; this gives a $\Sigma^{\frac{1}{2}}$ of 0.4 km corresponding to the MPE_2 column in the table.

Some comments may be useful on the reason for some of the larger sigmas in Table 2. The large $\sigma(M_1)$ (and hence $\sigma(a)$) and $\sigma(M_2)$ at the first epoch arose partly because this epoch was only 8 hours after launch and partly because of the complete absence of observations between $2^h 43^m$ on 6 May 1967 and $20^h 18^m$ on 7 May; there was a correlation of -0.988 between the computed values of M_1 and M_2 . Similarly, a large $\sigma(M_1)$ arose for epoch 1967 DEC 5 because of the missing data for 3-4 December, which has already been mentioned. High sigmas for epochs from 1968 DEC 20 to 1969 JAN 16, inclusive, arose because of the paucity of observations during this period; with only 14 observations accepted, the run of 1969 JAN 1 gave the largest sigmas, for the parameters Ω and ω , of all the runs in Table 2. High sigmas for epochs from 1969 FEB 27 to 1969 APR 4 arose partly from paucity of observations and partly from the large values of ϵ during that period; the epochs 1969 MAR 20 and 1969 MAR 23, for which the largest values of ϵ (and of sigmas for e , i and $M_0 + \omega$) of all the runs in Table 2 were obtained, have already been discussed in section 4.

6 DISCUSSION

Although the perigee height of Ariel 3, during the period considered, was around 500 km, as opposed to about 300 km for Ariel 2, air drag was still important enough to be the chief limitation in the computation of orbital parameters by PROP. The value of M_2 , equal to half the mean angular acceleration of the satellite, was only about $0.02^\circ/d^2$ immediately after launch; when this parameter approached or exceeded $0.1^\circ/d^2$, as it did for a few days in October-November 1968 and for longer periods in 1969, or when it changed by more than about $0.1^\circ/d^2$ from epoch to epoch, for example in late December 1967, the orbit does not fit the data so well, as indicated by higher values of ϵ . The period of validity of a set of orbital parameters is the same as the period spanned by the observations used in determining the parameters, i.e. between 3 and 4 days. If a set of orbital parameters is used to predict beyond the period of validity, then, when M_2 is changing rapidly (and ϵ is large), error increases rapidly. Now since each set of orbital

parameters (after the first) was obtained by iteration from an initial set equivalent to the parameters at the preceding epoch, an immediate guide to the accuracy of three-day prediction - i.e. up to five days from a given epoch - is provided by the largest absolute value for the residuals in the first iteration of the orbit determination at the next epoch. This largest absolute value can change violently. As an extreme example, the figures for a series of successive epochs, starting at 1969 OCT 12, are:- 29, 45, 13, 7, 12, 67, 154, 189, 628, 52, 30, 59, 46, 17, 12, 74, 10; the value 628, equivalent to an angular error of about 10° , occurred for epoch 1968 NOV 5, and was obviously due to the unusually high value (at 1968 NOV 2) of 0.1144 for M_2 , which immediately afterwards fell to 0.0530. (A very large magnetic storm occurred on 1 November and had a devastating effect on the upper atmosphere⁹.)

The behaviour of the orbital inclination, as evidenced by Fig.1, is worth discussing in detail. There are two distinct features. First, and very striking, is the secular behaviour: i remained essentially constant at 80.18° until the middle of December 1967, then dropped to about 80.163 in a period of about three months, and thereafter again remained essentially constant. Second, there are the superimposed oscillations, in which certain frequencies and amplitudes can fairly readily be seen. It is not entirely easy to explain either of these features.

Apart from resonances - and there should be no relevant resonance associated with the orbital parameters of Ariel 3 - the only known cause of secular variation in the orbital inclination of an earth satellite is the rotation of the atmosphere. Applying the formula of King-Hele and Scott¹⁰, if the atmosphere at a height of 500 km is taken to rotate at twice the angular velocity of the earth ($\Lambda = 2.0$ in Ref.10), then i in Fig.1 should show a secular drop of about 0.02° , i.e. just about what it does show. The rate of drop should be proportional to M_1 , however, whereas in Fig.1, as already remarked, the total drop is concentrated into a period of about three months, starting in December 1967.

The phenomenon is sufficiently remarkable for a sceptical reader to wonder whether the inclinations in Fig.1 really are right. Here the SDC values, though less accurate than the R.A.E. values, are useful; they are completely independent of the R.A.E. values, and confirm - not that there was a serious doubt - the secular behaviour indicated by the R.A.E. values.

Slightly more credible, though still unlikely, is the possibility that the sharp drop in i is not a purely secular effect but an oscillation superimposed on the change due to atmospheric rotation. Such an oscillation would require a period of 500 days or more, however, and even then the next cycle might have been expected to appear before the end of the graph. The amplitude of the oscillation would have to be nearly 0.01° . The author is unable to see whence such a term could arise.

There remain two possibilities: a single, complete discontinuity, due for example to meteoric impact, and a genuine (secular) perturbation over the (roughly) three-month period. The former seems very unlikely, though a discontinuity near the beginning of January 1968 cannot be completely ruled out; so we are left with the possibility of a real perturbation. Bearing in mind that a perturbing force, to produce an effect on i without affecting i , has to act in a direction perpendicular to the orbital plane, and that, to avoid cancellation, it has to act in opposite directions at the ascending and descending nodes of the orbit, it is difficult to see what the force can be, other than atmospheric rotation.

Attempts to produce an explanation must therefore degenerate into mere speculation. The three-month period of interest corresponded to a period of maximum solar activity (mean 107 mm solar radiation in excess of $150 \times 10^{-22} \text{ W m}^{-2} \text{ Hz}^{-1}$), and during this period M_2 was greater than before and after, though not enough greater to explain Fig.1 at once. Could it be that, at heights above 400 km, where no accurate measurements have been made, atmospheric rotation is significantly faster during periods of high solar activity, i.e. that King-Hele and Scott's Λ parameter is considerably less than 2.0 for most of Fig.1, but very much larger during the short period of maximum activity?

Apart from the correlation with solar activity, two other interesting (and unexplained) correlations should be mentioned. First, the direction of the spin axis of Ariel 3 has been monitored by RSRS, Slough¹. At injection the spin axis pointed 69° south, i.e. to a point on the celestial sphere with declination -69° . During the first three months the axis looped towards the south, reaching declination -86° on 19 June 1967, but after this it moved north and the declination remained positive after August 1967, for as long as the satellite was still spinning. The axis looped towards the north, reaching a declination of almost 90° on 25 or 26 February 1968; from early

January to early April 1968, i.e. roughly the period of sharp fall in orbital inclination, was the period during which the declination of the spin axis exceeded 45° . (There was another period, starting about the middle of September 1968, when the declination again exceeded 45° , but this only lasted for about one month instead of three.) Second, the satellite's tape recorder, which contained two litres of air at a pressure of one atmosphere, was operating only intermittently during the period of interest. Until 28 November 1967 the tape recorder worked successfully. It then failed, but recovered and worked, apparently perfectly again, for three periods (of two, four and three weeks), until it failed for the last time on 14 April 1968. It is tempting to speculate that air was leaking, but, even if this was so, it is difficult to see how the right inclination-reducing force could result.

On turning to the oscillations, it is clear from Fig.1 that a number of components, of differing frequency, amplitude and phase, are present. Since $\cos i$ is small and e is very small, the oscillatory perturbations due to the earth's odd harmonics, the amplitude of which is proportional to $e \cos i$, is completely negligible. The effect of the earth's tesseral harmonics, as remarked in section 3, is not negligible, and the amplitude of the $J_{2,2}$ perturbation is more than twice as big as some of the values of $\sigma(i)$ in Table 2; however, such effects should not appear in Fig.1, since they are averaged out during orbit determination. Hence the oscillatory components in Fig.1 may be thought of as being due solely to luni-solar perturbations, for which the various terms in di/dt are given in Ref.11 (equation (31)). The main term in the integral of the equation for di/dt is, for Ariel 3, $0^\circ.0015 \cos 2(u_s - \Omega)$, where u_s is the argument of latitude of the sun; the period of this term is 80 days and a complete cycle may be seen, in Fig.1, for example between MJD 39673 and 39753 and between MJD 39993 and 40073. The next largest terms are combined terms for the sun and moon which, if we ignore the small non-zero value (Ω_m) for the right ascension of the node of the moon's orbit, are given by $0^\circ.0012 \cos \Omega$ and $0^\circ.0007 \cos 2\Omega$, of period 280 days and 140 days respectively. The fourth largest term is also the principal one in which u_m , the argument of latitude of the moon, appears; ignoring Ω_m again, it is $0^\circ.0005 \cos 2(u_m - \Omega)$, of period 12½ days approximately. Other terms are of smaller argument, but a combination of such terms could produce a detectable contribution to the graph of i .

In the absence of a spectral analysis or a complete analysis of all terms from Ref.11 it is difficult to be sure whether the oscillatory component of Fig.1 can be fully explained by lumi-solar perturbations. It does appear, however, that Fig.1 contains a sinusoidal term with period about 30 days. A term in $\cos(k\Omega \pm u_m)$, for small integral k , would be appropriate here, but the only terms which arise have the eccentricity of the moon, i.e. 0.055, as a factor, and their amplitudes are too small.

One other known source of sinusoidal contributions to inclination variation should be mentioned. This is the precession and nutation of the earth's axis, which provides the reference with respect to which (the complement of) an orbital inclination is measured. The main contribution is from precession and may be taken from Ref.12 (in which the nutation terms have the wrong sign but the precession term is correct). For Ariel 3 this gives $-0.0007 \cos \Omega$, and so reduces (to about half) the amplitude of the direct lumi-solar perturbation term of argument Ω .

7 CONCLUSIONS

Orbital parameters for the satellite Ariel 3, as for Ariel 2, have been determined at R.A.E., at three-day intervals, from Minitrack observations supplied by NASA. The accuracy of the computed parameters is about the same, in general, as was obtained for Ariel 2, i.e. better than 1 km except on very rare occasions, and should be adequate for experimenters' requirements. (Values of eccentricity were accurate enough to be used in determining the odd harmonics in the geopotential¹³.) During periods of high drag, better accuracy could have been obtained by inclusion of an eighth parameter, M_3 , in the orbital model, as was done with Ariel 2.

Orbital inclination was determined rather less accurately than for Ariel 2, no doubt due to the fact that there are no Minitrack stations, in either hemisphere, at latitudes as high as 80° . The accuracy was good enough, however, for an anomalous secular behaviour in inclination to be clearly observable. This behaviour has been discussed but not explained.

Acknowledgement

The author wishes to thank Jennifer Davies for preparing the card decks and supervising the computer analysis.

Table 1
MINITRACK STATIONS OBSERVING ARIEL 3

Station name	Location	x (km)	y (km)	z (km)
Fort Myers	Fort Myers, Florida, U.S.A.	807.885	-5652.020	2833.549
Johannesburg	Hartebeshoek, South Africa	5084.798	2670.474	-2768.164
Lima	Lima, Peru	1388.818	-6088.429	-1293.207
Newfoundland	St. Johns, Newfoundland	2602.801	-3419.184	4697.694
Quito	Quito, Ecuador	1263.617	-6255.010	-68.856
Santiago	Santiago, Chile	1769.707	-5044.642	-3468.192
Winkfield	Winkfield, England	3983.130	-48.404	4964.711
Ulaska	Fairbanks, Alaska	-2282.332	-1452.667	5756.942
Madagascar	Tananarive, Malagasy	4091.903	4434.373	-2064.537
Orroral	Canberra, Australia	-4447.361	2677.215	-3695.209

Table 2

DATE	HJD	E	MJD
MAY 9 967	501.18	501.70	39616
MAY 12 967	503.00	503.00	39619
MAY 15 967	504.56	504.56	39622
MAY 18 967	506.08	506.08	39625
MAY 21 967	508.96	508.96	39628
MAY 24 967	508.73	508.73	39631
MAY 27 967	509.08	509.08	39634
MAY 30 967	509.92	509.92	39637
JUN 2 967	510.01	510.01	39640
JUN 5 967	509.50	509.50	39643
JUN 8 967	508.42	508.42	39646
JUN 11 967	506.93	506.93	39649
JUN 14 967	509.7	509.7	39652
JUN 17 967	509.16	509.16	39655
JUN 20 967	511.16	511.16	39658
JUN 23 967	513.16	513.16	39661
JUN 26 967	515.16	515.16	39664
JUN 29 967	517.16	517.16	39667
JUL 2 967	519.16	519.16	39670
JUL 5 967	521.16	521.16	39673
JUL 8 967	523.16	523.16	39676
JUL 11 967	525.16	525.16	39679
JUL 14 967	527.16	527.16	39682
JUL 17 967	529.16	529.16	39685
JUL 20 967	531.16	531.16	39688
JUL 23 967	533.16	533.16	39691
JUL 26 967	535.16	535.16	39694
JUL 29 967	537.16	537.16	39697
AUG 1 967	539.16	539.16	39700
AUG 4 967	540.56	540.56	39703
AUG 7 967	542.56	542.56	39706
AUG 10 967	544.56	544.56	39709
AUG 13 967	546.56	546.56	39712
AUG 16 967	548.56	548.56	39715
AUG 19 967	550.56	550.56	39718
AUG 22 967	552.56	552.56	39721
AUG 25 967	554.56	554.56	39724
AUG 28 967	556.56	556.56	39727
AUG 31 967	558.56	558.56	39730

Table 2. Orbital parameters of Ariel 3

Table 2(cont'd.)

DATE	a	e	h_p	i	Ω	ω	$H_0 + \omega$	H_1	H_2	D	E	HJD	
1967 SEP 3	6925.1506	5	0.007477	6	500.89	4	80.1831	9	16.8056	6	142.00	5	5626.7539
SEP 6	6925.1570	7	0.007968	6	502.76	4	80.1817	13	12.9895	7	153.20	4	5626.9814
SEP 9	6924.7797	4	0.008669	5	504.43	4	80.1808	14	9.1667	5	124.74	5	5625.1899
SEP 12	6926.8016	5	0.008150	6	505.94	4	80.1819	16	116.41	6	505.117	6	5625.3949
SEP 15	6924.6180	5	0.008207	6	507.17	4	80.1831	12	1.5257	8	108.00	7	5625.6132
SEP 18	6926.6613	5	0.008248	6	507.92	4	80.1849	11	357.7073	11	99.67	7	5625.7995
SEP 21	6924.2779	7	0.008281	10	507.99	7	80.1839	12	353.8861	15	91.43	10	5626.0151
SEP 24	6924.1083	6	0.008284	9	507.57	6	80.1819	12	350.0328	12	83.31	6	5626.7144
SEP 27	6925.9703	7	0.008283	11	506.55	7	80.1823	14	346.2698	14	75.32	6	5626.3768
SEP 30	6923.7676	8	0.008232	10	505.12	7	80.1824	14	342.4176	15	66.89	13	5626.6192
OCT 3	6923.6107	5	0.008151	7	503.40	5	80.1817	10	338.5932	11	58.43	7	5626.7997
OCT 6	6923.4613	4	0.008037	6	501.53	4	80.1793	8	334.7685	9	49.74	5	5626.9158
OCT 9	6923.3657	5	0.007903	5	499.63	5	80.1786	14	330.9434	11	40.97	6	5627.1112
OCT 12	6923.7119	4	0.007719	7	498.10	5	80.1777	11	327.1172	8	32.01	4	5627.3157
OCT 15	6923.0109	4	0.007550	7	496.82	5	80.1782	13	323.2937	6	22.94	4	5627.5051
OCT 18	6922.8410	4	0.007354	8	496.33	5	80.1783	9	4679.6	13	17.77	5	5627.902
OCT 21	6922.6723	5	0.007164	10	496.72	7	80.1781	12	315.6626	9	6.42	2	5627.9035
OCT 24	6922.4796	5	0.006900	10	498.19	7	80.1791	11	311.8163	10	354.89	5	5627.9352
OCT 27	6922.2590	9	0.006679	18	500.57	12	80.1774	19	307.9862	19	345.25	10	5628.3898
OCT 30	6921.0522	9	0.006472	18	503.64	12	80.1765	18	304.1576	17	18.97	10	5628.5068
NOV 2	6921.6745	8	0.006267	15	507.64	10	80.1753	16	300.3293	15	325.02	8	5629.0774
NOV 5	6921.4201	5	0.006072	9	511.80	6	80.1763	11	296.4997	10	314.44	6	5629.3770
NOV 8	6921.2304	5	0.005914	9	515.94	6	80.1766	12	292.6709	11	303.62	7	5629.9327
NOV 11	6921.0603	5	0.005802	7	519.37	5	80.1772	13	288.8397	10	292.52	8	5629.8007
NOV 14	6920.8900	4	0.005731	6	521.62	4	80.1769	9	285.6093	9	281.26	6	5630.0011
NOV 17	6920.7058	4	0.005709	5	522.26	4	80.1775	11	281.1777	8	269.70	8	5630.2180
NOV 20	6920.4900	6	0.005710	7	521.23	5	80.1781	14	277.3462	11	258.24	9	5630.4614
NOV 23	6920.2475	7	0.005793	11	518.51	8	80.1769	19	273.5164	16	266.94	11	5630.7577
NOV 26	6919.9821	5	0.005680	8	514.73	5	80.1761	15	269.6824	10	235.65	8	5630.0260
NOV 29	6919.7398	7	0.005621	9	510.16	6	80.1789	16	265.8547	10	224.52	7	5630.9586
DEC 2	6919.5008	7	0.006155	10	505.79	7	80.1789	17	262.0189	10	213.80	7	5631.6136
DEC 5	6919.3623	19	0.006366	8	501.49	6	80.1781	13	258.1892	8	203.40	6	5631.8240
DEC 8	6919.1521	5	0.006557	9	498.10	6	80.1790	13	254.3560	7	193.21	5	5632.0481
DEC 11	6918.9895	5	0.006725	9	495.92	6	80.1804	12	250.5260	7	183.67	5	5632.5169
DEC 14	6918.8397	6	0.006913	11	494.62	7	80.1814	14	246.6908	10	173.70	6	5632.1611
DEC 17	6918.6310	6	0.007125	11	494.01	8	80.1813	16	242.8597	13	164.46	7	5632.6620
DEC 20	6918.5194	8	0.007327	15	494.13	10	80.1802	19	239.0276	19	155.22	10	5633.0159
DEC 23	6918.0284	7	0.007546	14	494.63	10	80.1789	15	235.1930	18	146.42	9	5633.5722
DEC 26	6917.4217	5	0.007691	9	496.03	6	80.1791	10	231.3581	11	137.61	5	5633.8424
DEC 29	6917.6294	7	0.007723	13	497.58	9	80.1788	16	227.5230	15	128.99	8	5633.8424

* FOR THESE TWO EPOCHS THE ADDITIONAL PARAMETER M_3 WAS INCLUDED IN THE ORBITAL MODEL. THE VALUE OBTAINED, IN BOTH CASES, WAS -0.00069 3.

Table 2(cont'd.) Orbital parameters of Ariel 3

Table 2(cont'd.)

DATE	a	e	h_p	ι	Ω	ω	$h_0 + \omega$	H_1	H_2	N	K	D	E	MJD	
1968 JAN 1	6917.4121	7	0.0079699	12	498.75	80.1787	14	223.6878	15	120.67	9	4.4	2.5	39856	
JAN 4	6917.1810	5	0.008167	11	499.94	80.1756	10	216.8503	10	112.14	6	3.6	2.1	39859	
JAN 7	6916.9663	5	0.0081167	9	500.68	80.1759	10	216.0118	11	103.74	6	3.6	2.2	39862	
JAN 10	6916.7298	7	0.008191	12	500.98	80.1743	15	212.1707	15	95.49	8	3.6	2.2	39863	
JAN 13	6916.4463	8	0.008208	10	500.70	7	80.1733	15	208.3333	15	87.04	7	3.6	2.2	39865
JAN 16	6916.1630	5	0.008186	15	499.55	10	80.1721	21	204.4944	14	78.98	5	3.6	2.2	39866
JAN 19	6915.9404	5	0.008166	8	498.57	5	80.1709	11	200.6511	10	70.47	6	3.6	2.2	39868
JAN 22	6915.7573	5	0.008104	8	496.92	6	80.1731	11	196.8099	10	62.13	7	3.6	2.2	39870
JAN 25	6915.5907	5	0.008019	9	495.00	6	80.1728	12	192.9679	15	53.69	8	3.6	2.2	39872
JAN 28	6915.3837	7	0.007612	9	492.92	6	80.1697	15	189.1269	15	45.03	9	3.6	2.2	39874
JAN 31	6915.1114	5	0.007789	7	490.85	5	80.1701	13	185.2848	12	36.35	8	3.6	2.2	39876
FEB 3	6914.7760	7	0.007657	9	487.01	7	80.1697	15	181.4404	15	27.54	10	3.6	2.2	39878
FEB 6	6914.4047	8	0.007487	12	487.84	8	80.1698	17	177.5952	16	18.59	10	3.6	2.2	39880
FEB 9	6914.1237	12	0.007302	20	487.50	14	80.1680	29	173.7501	25	9.32	15	3.6	2.2	39882
FEB 12	6913.8176	12	0.007016	18	488.09	13	80.1698	24	169.9029	24	55.97	17	3.6	2.2	39884
FEB 15	6913.6033	10	0.006955	18	489.37	18	80.1704	23	166.6032	22	250.85	14	3.6	2.2	39886
FEB 18	6913.4091	6	0.006768	11	491.89	8	80.1684	14	162.2163	13	340.99	6	3.6	2.2	39888
FEB 21	6913.1948	5	0.006558	11	495.41	8	80.1651	10	158.3698	10	330.99	6	3.6	2.2	39890
FEB 24	6912.9728	5	0.006398	10	499.24	7	80.1651	11	154.5200	11	320.67	6	3.6	2.2	39892
FEB 27	6912.7332	6	0.006235	12	503.38	8	80.1647	12	150.6704	12	310.07	8	3.6	2.2	39894
MAR 1	6912.4088	7	0.006099	12	507.11	8	80.1633	14	146.8197	13	299.21	10	3.6	2.2	39896
MAR 4	6912.1074	6	0.006055	10	510.13	7	80.1616	12	142.9574	12	287.93	9	3.6	2.2	39898
MAR 7	6911.8210	7	0.005909	11	511.82	8	80.1612	12	140.1139	13	276.41	12	3.6	2.2	39900
MAR 10	6911.6263	5	0.005882	7	511.90	9	80.1610	11	135.2602	11	265.01	10	3.6	2.2	39902
MAR 13	6911.4428	5	0.005899	7	510.30	5	80.1603	11	131.4055	11	253.48	9	3.6	2.2	39904
MAR 16	6911.2472	5	0.005966	7	507.15	5	80.1605	11	127.5526	12	241.80	9	3.6	2.2	39906
MAR 19	6911.0661	4	0.005669	7	503.05	5	80.1583	12	123.6953	11	230.05	7	3.6	2.2	39908
MAR 22	6910.8834	4	0.006201	7	498.42	5	80.1596	9	119.8392	8	219.67	5	3.6	2.2	39910
MAR 25	6910.6620	4	0.006364	7	494.24	5	80.1601	11	115.9846	6	209.02	4	3.6	2.2	39912
MAR 28	6910.3900	4	0.006526	7	499.49	5	80.1584	10	112.1245	6	198.49	4	3.6	2.2	39914
MAR 31	6910.0864	5	0.006751	5	487.38	7	80.1589	11	108.2677	6	185.70	4	3.6	2.2	39916
APR 3	6909.7825	4	0.006935	7	485.26	5	80.1606	15	104.4090	6	179.16	4	3.6	2.2	39918
APR 6	6908.5307	6	0.007131	9	484.21	6	80.1615	15	103.5330	10	169.88	6	3.6	2.2	39920
APR 9	6909.3229	6	0.007302	6	484.14	6	80.1630	11	96.4947	5	160.92	5	3.6	2.2	39922
APR 12	6909.1514	5	0.007459	9	484.87	6	80.1653	15	92.8358	12	152.18	7	3.6	2.2	39924
APR 15	6908.9427	4	0.007406	5	483.07	6	80.1661	9	88.9791	8	143.42	4	3.6	2.2	39926
APR 18	6908.7402	4	0.007727	7	487.66	5	80.1642	12	85.1292	9	134.77	5	3.6	2.2	39928
APR 21	6908.5658	3	0.007824	4	489.41	3	80.1661	7	81.2430	6	126.28	3	3.6	2.2	39930
APR 24	6908.3844	4	0.007837	6	490.79	4	80.1643	10	77.4039	9	117.73	3	3.6	2.2	39932
APR 27	6908.1869	4	0.008026	6	491.87	4	80.1663	10	73.5462	10	109.23	5	3.6	2.2	39934

Table 2(cont'd.) Orbital parameters of Ariel 3

Table 2(cont'd.)

DATE	a	e	h _p	i	Ω	ω + ω	H ₂	H ₁	E	MJD
1968 APR 30	6907.9801	3	0.008728	7	492.61	5	60.1625	11	65.0286	12
MAY 1	6907.7640	3	0.008731	7	492.62	5	60.1625	10	61.0656	10
MAY 2	6907.5427	5	0.008732	7	492.27	5	60.1621	10	61.0656	10
MAY 3	6907.3210	5	0.008715	7	491.27	5	60.1628	9	58.1033	11
MAY 4	6906.9929	5	0.008734	6	490.10	6	60.1641	10	54.2407	9
MAY 5	6906.7581	6	0.007725	6	482.58	6	60.1654	10	50.3793	6
MAY 6	6906.5194	3	0.007737	6	482.88	4	60.1661	11	46.9160	5
MAY 7	6906.2481	3	0.007745	5	482.15	4	60.1666	10	42.6350	6
MAY 8	6905.9706	4	0.007719	6	482.55	4	60.1682	10	38.7904	7
MAY 9	6905.7282	4	0.007729	7	482.31	5	60.1690	10	34.9287	8
MAY 10	6905.5046	3	0.007704	7	481.73	5	60.1687	11	31.0642	4
MAY 11	6905.3086	3	0.006816	6	481.82	5	60.1663	7	27.2011	7
MAY 12	6905.1591	4	0.007645	5	482.88	5	60.1649	9	23.3359	9
MAY 13	6905.0339	4	0.006642	3	482.87	5	60.1651	9	19.4705	7
MAY 14	6904.8462	4	0.006621	5	482.00	5	60.1626	14	15.6051	6
MAY 15	6904.6736	3	0.006616	6	481.83	4	60.1628	7	11.7337	7
MAY 16	6904.5393	3	0.005869	5	492.20	5	60.1642	9	7.8672	6
MAY 17	6904.4037	3	0.005769	4	50.29	5	60.1641	8	5.9988	5
MAY 18	6904.2859	4	0.005690	4	50.53	5	60.1648	8	30.3.82	5
MAY 19	6904.1790	3	0.005659	4	50.57	5	60.1649	8	29.2.24	5
MAY 20	6904.0683	3	0.005641	4	50.63	5	60.1673	7	28.0.94	5
MAY 21	6903.9553	4	0.005686	4	50.63	5	60.1642	9	35.6.26	5
MAY 22	6903.8568	4	0.005686	4	50.63	5	60.1669	10	34.8.9271	10
MAY 23	6903.7691	5	0.005715	5	50.54	4	60.1653	9	34.6.6593	9
MAY 24	6903.6649	6	0.005762	6	49.4.66	6	60.1637	9	36.0.7904	10
MAY 25	6903.5221	4	0.006622	2	49.0.46	5	60.1642	7	33.6.021	7
MAY 26	6903.4179	4	0.006622	2	49.0.46	5	60.1642	7	33.5.0530	7
MAY 27	6903.3169	3	0.006623	2	48.6.42	5	60.1641	11	32.9.1336	6
MAY 28	6903.2169	3	0.006623	2	48.3.42	5	60.1638	9	32.5.0672	6
MAY 29	6903.1169	3	0.006623	2	48.0.42	5	60.1635	8	30.5.6672	8
MAY 30	6903.0169	3	0.006623	2	47.7.42	5	60.1633	8	30.2.0661	9
MAY 31	6902.9169	3	0.006623	2	47.4.42	5	60.1633	8	29.8.2265	6
JUN 1	6902.8169	4	0.006623	2	47.1.42	5	60.1693	10	31.7.5742	9
JUN 2	6902.7169	4	0.006623	2	46.8.42	5	60.1642	7	31.5.0530	7
JUN 3	6902.6169	4	0.006623	2	46.5.42	5	60.1642	7	31.3.7061	9
JUN 4	6902.5169	4	0.006623	2	46.2.42	5	60.1641	11	30.9.8374	8
JUN 5	6902.4169	4	0.006623	2	45.9.42	5	60.1638	9	30.5.672	8
JUN 6	6902.3169	4	0.006623	2	45.6.42	5	60.1635	8	30.2.0661	9
JUN 7	6902.2169	4	0.006623	2	45.3.42	5	60.1643	9	29.8.2265	6
JUN 8	6902.1169	4	0.006623	2	45.0.42	5	60.1637	9	29.8.1627	7
JUN 9	6902.0169	4	0.006623	2	44.7.42	5	60.1631	11	29.4.3529	7
JUN 10	6901.9169	4	0.006623	2	44.4.42	5	60.1631	11	29.0.4617	9
JUN 11	6901.8169	4	0.006623	2	44.1.42	5	60.1631	11	28.6.6073	9
JUN 12	6901.7169	4	0.006623	2	43.8.42	5	60.1631	11	28.2.7332	8
JUN 13	6901.6169	4	0.006623	2	43.5.42	5	60.1631	10	27.8.8640	9
JUN 14	6901.5169	4	0.006623	2	43.2.42	5	60.1631	10	27.5.8640	9
JUN 15	6901.4169	4	0.006623	2	42.9.42	5	60.1631	10	27.2.8640	9
JUN 16	6901.3169	4	0.006623	2	42.6.42	5	60.1631	10	26.9.8640	9
JUN 17	6901.2169	4	0.006623	2	42.3.42	5	60.1631	10	26.6.8640	9
JUN 18	6901.1169	4	0.006623	2	42.0.42	5	60.1631	10	26.3.8640	9
JUN 19	6901.0169	4	0.006623	2	41.7.42	5	60.1631	10	26.0.8640	9
JUN 20	6900.9169	4	0.006623	2	41.4.42	5	60.1631	10	25.7.8640	9
JUN 21	6900.8169	4	0.006623	2	41.1.42	5	60.1631	10	25.4.8640	9
JUN 22	6900.7169	4	0.006623	2	40.8.42	5	60.1631	10	25.1.8640	9
JUN 23	6900.6169	4	0.006623	2	40.5.42	5	60.1631	10	24.8.8640	9
JUN 24	6900.5169	4	0.006623	2	40.2.42	5	60.1631	10	24.5.8640	9
JUN 25	6900.4169	4	0.006623	2	39.9.42	5	60.1631	10	24.2.8640	9
JUN 26	6900.3169	4	0.006623	2	39.6.42	5	60.1631	10	23.9.8640	9
JUN 27	6900.2169	4	0.006623	2	39.3.42	5	60.1631	10	23.6.8640	9
JUN 28	6900.1169	4	0.006623	2	39.0.42	5	60.1631	10	23.3.8640	9
JUN 29	6900.0169	4	0.006623	2	38.7.42	5	60.1631	10	23.0.8640	9
JUN 30	6900.9169	4	0.006623	2	38.4.42	5	60.1631	10	22.7.8640	9
JUN 31	6900.8169	4	0.006623	2	38.1.42	5	60.1631	10	22.4.8640	9
JUL 1	6900.7169	4	0.006623	2	37.8.42	5	60.1631	10	22.1.8640	9
JUL 2	6900.6169	4	0.006623	2	37.5.42	5	60.1631	10	21.8.8640	9
JUL 3	6900.5169	4	0.006623	2	37.2.42	5	60.1631	10	21.5.8640	9
JUL 4	6900.4169	4	0.006623	2	36.9.42	5	60.1631	10	21.2.8640	9
JUL 5	6900.3169	4	0.006623	2	36.6.42	5	60.1631	10	20.9.8640	9
JUL 6	6900.2169	4	0.006623	2	36.3.42	5	60.1631	10	20.6.8640	9
JUL 7	6900.1169	4	0.006623	2	36.0.42	5	60.1631	10	20.3.8640	9
JUL 8	6900.0169	4	0.006623	2	35.7.42	5	60.1631	10	19.0.8640	9
JUL 9	6900.9169	4	0.006623	2	35.4.42	5	60.1631	10	18.7.8640	9
JUL 10	6900.8169	4	0.006623	2	35.1.42	5	60.1631	10	18.4.8640	9
JUL 11	6900.7169	4	0.006623	2	34.8.42	5	60.1631	10	18.1.8640	9
JUL 12	6900.6169	4	0.006623	2	34.5.42	5	60.1631	10	17.8.8640	9
JUL 13	6900.5169	4	0.006623	2	34.2.42	5	60.1631	10	17.5.8640	9
JUL 14	6900.4169	4	0.006623	2	33.9.42	5	60.1631	10	17.2.8640	9
JUL 15	6900.3169	4	0.006623	2	33.6.42	5	60.1631	10	16.9.8640	9
JUL 16	6900.2169	4	0.006623	2	33.3.42	5	60.1631	10	16.6.8640	9
JUL 17	6900.1169	4	0.006623	2	33.0.42	5	60.1631	10	16.3.8640	9
JUL 18	6900.0169	4	0.006623	2	32.7.42	5	60.1631	10	16.0.8640	9
JUL 19	6900.9169	4	0.006623	2	32.4.42	5	60.1631	10	15.7.8640	9
JUL 20	6900.8169	4	0.006623	2	32.1.42	5	60.1631	10	15.4.8640	9
JUL 21	6900.7169	4	0.006623	2	31.8.42	5	60.1631	10	15.1.8640	9
JUL 22	6900.6169	4	0.006623	2	31.5.42	5	60.1631	10	14.8.8640	9
JUL 23	6900.5169	4	0.006623	2	31.2.42	5	60.1631	10	14.5.8640	9
JUL 24	6900.4169	4	0.006623	2	30.9.42	5	60.1631	10	14.2.8640	9
JUL 25	6900.3169	4	0.006623	2	30.6.42	5	60.1631	10	13.9.8640	9
JUL 26	6900.2169	4	0.006623	2	30.3.42	5	60.1631	10	13.6.8640	9
JUL 27	6900.1169	4	0.006623	2	30.0.42	5	60.1631	10	13.3.8640	9
JUL 28	6900.0169	4	0.006623	2	29.7.42	5	60.1631	10	13.0.8640	9
JUL 29	6900.9169	4	0.006623	2	2					

Table 2(cont'd.)

DATE	a	e	i	Ω	ω	ω₀ + ω₀	H₁	H₂	E	HJD
							h _p	h _o	h _o + h _o	h _o
1968 AUG 28	6891.2363	5	484.33	3	80.1593	8	274.9904	7	5453.2227	40096
AUG 31	6891.1391	5	487.16	4	80.1595	11	271.1167	10	5453.3379	40099
SEP 3	6890.9737	5	485.69	3	80.1585	9	267.2414	6	5453.3349	40102
SEP 6	6890.7820	3	484.07	3	80.1596	9	265.3647	51	5453.3359	40105
SEP 9	6890.5445	5	482.46	5	80.1617	13	259.4891	9	5453.3364	40108
SEP 12	6890.2949	6	480.79	5	80.1637	11	255.6141	11	5454.0432	40111
SEP 15	6890.9888	5	479.27	5	80.1611	12	251.7396	12	5454.3392	40114
SEP 18	6890.7099	7	478.52	6	80.1593	13	247.8628	14	5454.7023	40117
SEP 21	6890.4432	4	478.23	5	80.1593	9	263.9842	9	5455.0331	40120
SEP 24	6890.1494	4	479.14	5	80.1621	11	240.1058	7	5455.3693	40123
SEP 27	6890.8431	4	480.99	5	80.1637	12	236.2269	7	5455.6093	40126
SEP 30	6890.5021	6	483.97	6	80.1651	12	232.3469	12	5456.4465	40129
OCT 3	6890.1135	6	485.56	10	80.1627	13	228.4721	12	5456.9277	40132
OCT 6	6890.7831	5	491.69	6	80.1629	11	224.5910	11	5457.3200	40135
OCT 9	6890.4690	5	495.81	6	80.1623	10	220.7095	11	5457.6224	40138
OCT 12	6890.1632	7	499.47	6	80.1634	16	216.8296	14	5458.1641	40141
OCT 15	6890.8286	6	501.82	6	80.1628	18	212.9478	13	5458.4534	40144
OCT 18	6890.5598	4	502.27	4	80.1641	11	209.0681	8	5458.7361	40147
OCT 21	6890.2740	6	501.23	6	80.1670	12	205.1833	13	5459.1123	40150
OCT 24	6890.0058	13	505.59	18	80.1653	27	201.5062	29	5459.4352	40153
OCT 27	6890.6671	17	509.29	10	80.1643	16	197.4232	15	5459.8732	40156
OCT 30	6890.2881	9	505.14	9	80.1634	15	193.5377	9	5460.2433	40159
OCT 18	6890.5598	4	505.23	5	80.1628	13	192.6177	17	5460.5036	40162
OCT 21	6890.2740	6	505.12	6	80.1670	12	205.1833	13	5460.9089	40165
OCT 24	6890.0058	13	505.59	18	80.1653	27	201.5062	29	5461.3756	40168
OCT 27	6890.6671	17	509.29	10	80.1643	16	197.4232	15	5461.6670	40171
OCT 30	6890.2881	9	505.14	9	80.1634	15	193.5377	9	5462.2943	40174
NOV 2	6890.7621	7	505.52	11	80.1641	11	189.6515	12	5462.6003	40177
NOV 5	6890.3693	6	505.76	10	80.1669	15	185.7665	13	5462.9461	40180
NOV 8	6890.1261	6	505.62	9	80.1677	9	181.8826	13	5463.3200	40183
NOV 11	6890.8495	5	506.02	9	80.1653	12	180.8947	11	5463.5533	40186
NOV 14	6890.5965	4	505.21	8	80.1643	16	174.1077	8	5464.1947	40189
NOV 17	6890.3399	4	506.59	10	80.1634	15	170.2159	16	5464.5514	40192
NOV 20	6890.0485	5	506.50	7	80.1633	11	164.3274	11	5464.8211	40195
NOV 23	6890.8017	4	506.70	9	80.1616	11	162.4388	12	5464.8065	40198
NOV 26	6890.5514	6	506.62	9	80.1628	9	177.55	6	5465.5372	40201
NOV 29	6890.2691	4	506.96	2	80.1579	13	154.6575	8	5466.0605	40204
DEC 2	6890.0138	4	507.02	5	80.1622	10	150.7466	9	5466.8730	40207
DEC 5	6890.7244	5	506.98	7	80.1608	12	164.8746	11	5466.1769	40210
DEC 8	6890.4845	6	507.63	9	80.1587	13	162.9364	12	5466.3211	40213
DEC 11	6890.2708	4	506.82	8	80.1553	17	139.0901	9	5466.4033	40215
DEC 14	6890.0563	5	507.35	11	80.1623	12	161.7468	12	5466.8162	40218
DEC 17	6890.8845	4	506.96	9	80.1622	10	151.2986	14	5466.3209	40221
DEC 20	6890.6933	11	507.07	6	80.1524	37	127.4255	16	5466.1349	40224
DEC 23	6890.4331	4	506.64	37	80.1572	16	123.5193	11	5466.0378	40225

Table 2(cont'd.) Orbital parameters of Ariel 3

Table 2(cont'd.)

DATE	G	h _p	Q	W	M ₂	K	D	C	E	MJD	
1968 DEC 26	6890.1374	9	0.006742 4.5	477.68 30	80.1604 15	119.6199 16	50.77 20	381.37 06	4.6	402116	
1969 DEC 29	6889.8561	8	0.006608 4.6	475.62 32	80.1584 16	115.7192 16	41.93 27	131.35 73	4.6	402119	
1969 JAN 1	6889.5904	10	0.005590 58	473.01 40	80.1554 37	111.8612 48	33.69 41	322.31 01	6.2	402222	
1969 JAN 4	6889.3649	7	0.006604 26	471.89 18	80.1579 22	107.9390 11	25.39 31	134.1453 42	5.6	402223	
1969 JAN 7	6889.1349	9	0.005202 26	471.07 18	80.1633 19	106.0360 13	15.66 42	366.7948 6.7	5.6	402226	
1969 JAN 10	6889.9023	7	0.006604 17	471.11 12	80.1673 15	106.0360 14	5.63 45	185.2542 4.2	5.6	402251	
1969 JAN 13	6888.7267	8	0.005855 19	471.85 13	80.1661 26	96.2158 41	356.86 51	14.48 46	4.9	402234	
1969 JAN 16	6888.5214	7	0.005651 11	471.06 7	80.1603 15	97.3565 9	34.66 74	209.32 94	3.4	402237	
1969 JAN 19	6888.2698	5	0.005552 9	476.92 6	80.1619 14	88.4618 9	356.51 20	45.03 16	1.9	402240	
1969 JAN 22	6888.0800	4	0.005426 7	480.73 5	80.1662 9	84.3596 8	325.62 6	4.0 0.0458	4.5	402243	
1969 JAN 25	6887.9095	8	0.005144 17	484.76 12	80.1601 26	80.0524 30	314.89 11	78.6350 14	10.0 0.0664	10	402246
1969 JAN 28	6887.6762	4	0.004972 7	489.13 5	80.1601 11	76.7421 6	303.42	7 276.48 94	7	402249	
1969 JAN 31	6887.4917	6	0.004898 6	492.44 4	80.1630 12	72.8382 7	291.51	8 115.10 33	6	402252	
1969 FEB 3	6887.2796	9	0.004820 14	494.63 10	80.1647 27	68.1398 22	279.28	18 154.34 03	1.9	402255	
1969 FEB 6	6887.0048	6	0.006735 9	495.07 6	80.1545 20	65.0393 18	267.7	1.9 1.0 0.6938	1.7	402258	
1969 FEB 9	6886.7448	5	0.004800 4	493.58 3	80.1603 10	61.1389 5	255.14	7 359.46 19	5	402261	
1969 FEB 12	6886.4480	7	0.004698 9	490.21 6	80.1623 14	57.2259 12	243.30	10 197.72 23	9	402264	
1969 FEB 15	6886.2096	4	0.004973 5	486.14 3	80.1599 8	53.3220 7	231.36	7 40.86 75	6	402267	
1969 FEB 18	6886.0037	6	0.005041 11	481.54 8	80.1604 17	69.0164 12	219.73	11 244.79 57	1.9	402270	
1969 FEB 21	6885.7394	5	0.005277 9	476.90 6	80.1612 13	45.0080 7	208.79	3 89.37 63	6	402273	
1969 FEB 24	6885.3901	4	0.005671 8	472.86 6	80.1639 10	41.0039 8	198.27	6 293.19 13	7	402276	
1969 FEB 27	6885.0175	10	0.005662 16	469.78 11	80.1614 20	37.6972 21	188.05	23 142.23 37	2.6	402279	
1969 MAR 2	6884.5708	9	0.005868 16	467.75 11	80.1661 23	35.7915 20	177.72	14 350.79 06	2.6	402282	
1969 MAR 5	6884.2249	8	0.006078 33	466.54 22	80.1635 23	29.8816 21	168.13	10 200.81 74	1.9	402285	
1969 MAR 8	6883.8985	7	0.006259 24	466.51 16	80.1620 12	25.763 11	158.86	29 51.97 16	3.4	402288	
1969 MAR 11	6883.5562	6	0.006329 41	467.13 28	80.1648 21	22.0688 18	149.94	32 264.34 10	5.3	402291	
1969 MAR 14	6883.2018	7	0.006626 47	467.91 32	80.1645 19	18.1605 15	141.24	23 269.75 56	4.8	402294	
1969 MAR 17	6882.8018	10	0.006791 50	469.04 34	80.1637 19	16.2503 18	132.56	23 338.87 72	6.9	402297	
1969 MAR 20	6882.2606	13	0.006876 72	470.35 49	80.1664 32	10.3406 31	124.43	24 189.44 65	7.4	403000	
1969 MAR 23	6881.5632	15	0.006226 66	472.25 45	80.1682 43	6.6266 32	116.00	22 4.8 237.8	6.6	403003	
1969 MAR 26	6880.8331	11	0.007002 43	471.99 29	80.1629 25	2.5198 24	108.1	13 547.50 61	1.9	403006	
1969 MAR 29	6880.3043	9	0.007035 49	471.91 34	80.1670 23	358.6061 17	100.1	13 133.44 76	1.8	403009	
1969 APR 1	6879.7769	8	0.007015 22	472.55 15	80.1668 19	354.6895 15	92.4	13 359.02 23	2.5	403112	
1969 APR 4	6879.2599	9	0.006955 24	472.26 17	80.1622 25	350.77 47	85.80	21 226.45 04	2.7	403115	
1969 APR 7	6878.8026	13	0.006921 13	471.21 9	80.1640 14	346.8569 15	75.94	14 95.64 64	2.0	403118	
1969 APR 10	6878.3687	7	0.006795 10	470.01 7	80.1660 17	362.9124 16	67.49	13 326.42 19	1.5	403121	
1969 APR 13	6877.9565	10	0.006681 12	468.48 8	80.1627 15	339.0175 16	59.46	24 198.71 06	2.3	403124	
1969 APR 16	6877.4873	5	0.006554 17	466.49 5	80.1630 15	335.0973 11	51.05	13 72.93 88	1.5	403127	
1969 APR 19	6877.0112	7	0.006389 10	464.46 7	80.1611 19	331.798 20	42.28	20 308.13 63	1.6	403130	
1969 APR 22	6876.6664	5	0.006195 11	462.80 7	80.1607 15	327.2526 14	33.40	18 33.40 18	2.1	403133	

Table 2(cont'd.) Orbital parameters of Ariel 3

Table 2(concl'd.)

DATE	a	e	h_p	\dot{e}	Ω	ω	$\Omega + \omega$	H_1	H_2	K	D	E	MJD	
1969 APR 25	6.876.2819	5	0.006014	9	461.36	6	80.1628	12	323.3265	10	24.47	9	65.5871	9
APR 28	6.875.9602	5	0.005825	8	460.48	6	80.1614	12	319.4033	11	14.93	9	303.1928	9
MAY 1	6.875.6463	5	0.005616	7	460.62	5	80.1590	9	315.4794	7	5.86	5	183.9548	6
MAY 4	6.875.3863	4	0.005412	7	461.68	4	80.1623	10	311.5449	7	355.75	5	65.3778	12
MAY 7	6.875.1606	7	0.005219	12	463.77	5	80.1603	15	307.6229	14	365.59	10	308.3819	12
MAY 10	6.874.9046	4	0.005047	6	466.70	4	80.1612	11	303.6977	7	335.24	7	191.8551	6
MAY 13	6.874.5963	10	0.004851	14	470.78	10	80.1636	22	299.7671	23	323.86	24	76.3603	19
MAY 16	6.874.1685	11	0.004711	16	474.81	11	80.1597	21	295.8419	21	312.39	32	5484.9672	12
MAY 19	6.873.8195	7	0.004660	7	478.54	5	80.1605	14	291.9102	11	301.09	16	209.4250	11
MAY 22	6.873.5133	10	0.004528	9	481.52	6	80.1615	19	287.9846	15	289.23	18	97.8301	13
MAY 25	6.873.2236	4	0.004478	8	483.24	5	80.1597	9	284.0519	8	276.74	13	347.3093	10
MAY 28	6.872.9662	6	0.004494	11	482.96	7	80.1622	10	280.1204	10	266.41	11	237.7840	10
MAY 31	6.872.7293	5	0.004512	11	481.16	7	80.1617	9	276.1912	10	252.43	11	120.4257	10
JUN 3	6.872.5101	5	0.004593	12	487.70	8	80.1615	12	272.2571	9	240.47	8	21.2948	10
JUN 6	6.872.2487	6	0.004680	16	473.52	11	80.1618	10	268.3288	13	223.93	15	274.2790	17
JUN 9	6.871.8698	6	0.004795	15	468.85	11	80.1593	14	264.3932	9	217.45	9	168.3830	13
JUN 12	6.871.4237	6	0.004971	16	464.24	11	80.1623	11	260.4588	14	206.72	11	44.3884	14
JUN 15	6.870.9359	5	0.005120	14	460.51	10	80.1613	11	256.5276	9	196.07	9	321.3446	12
JUN 18	6.870.4774	13	0.005321	23	457.54	16	80.1632	26	252.5907	24	186.39	24	220.4339	23
JUN 21	6.870.1597	6	0.005473	15	456.04	11	80.1597	12	248.6569	15	176.19	15	120.8883	18
JUN 24	6.869.9059	6	0.005645	10	455.47	7	80.1593	13	244.7181	11	166.67	14	22.3648	12
JUN 27	6.869.6849	6	0.005835	12	455.58	8	80.1623	12	240.7828	14	157.56	14	284.6949	15
JUN 30	6.869.5072	4	0.005933	8	456.63	6	80.1584	15	236.8446	10	148.51	15	187.7126	13
JUL 3	6.869.2841	6	0.006124	8	457.98	6	80.1593	15	232.9084	20	139.79	16	91.4356	13
JUL 6	6.869.0454	5	0.006253	6	459.47	4	80.1597	10	228.9696	8	131.24	6	5491.3312	7
JUL 9	6.868.7656	7	0.006149	8	461.03	5	80.1595	11	225.9307	14	122.80	13	261.4793	13
JUL 12	6.868.4837	5	0.006442	6	462.27	4	80.1572	11	221.0918	8	114.44	8	167.9902	8
JUL 15	6.868.2055	5	0.006472	8	463.39	5	80.1596	12	217.1502	11	106.46	7	66.4469	13
JUL 18	6.868.0054	4	0.006532	6	463.82	4	80.1574	8	213.2111	8	98.37	6	143.9226	7
JUL 21	6.867.8334	5	0.006515	5	466.13	3	80.1588	10	209.2683	9	90.44	6	5492.9781	6
JUL 24	6.867.6693	4	0.006514	5	465.66	3	80.1578	10	205.3303	8	82.23	6	162.3519	7
JUL 27	6.867.4870	3	0.006444	6	462.98	4	80.1560	12	201.3852	8	74.26	7	72.8838	7
JUL 30	6.867.2866	7	0.006365	9	461.73	6	80.1578	14	197.4469	13	66.10	12	343.5886	11
AUG 2	6.867.0288	6	0.006254	7	460.22	5	80.1590	12	193.5018	10	58.14	12	255.0602	11
AUG 5	6.866.6660	6	0.006114	9	458.40	6	80.1591	12	189.5630	13	49.86	17	168.5558	14
AUG 8	6.866.2557	5	0.005933	7	456.60	5	80.1606	12	185.6171	10	41.28	14	82.7318	13
AUG 11	6.865.8483	6	0.005744	8	454.97	5	80.1594	11	181.6753	11	32.73	15	358.3244	12
AUG 14	6.865.4953	7	0.005556	13	453.71	9	80.1603	14	177.7275	16	24.16	20	275.6700	17
AUG 17	6.865.2333	5	0.005317	7	453.29	5	80.1625	11	173.7831	9	14.83	18	195.9412	14
AUG 20	6.865.0070	10	0.005100	28	453.57	19	80.1600	15	169.8373	16	5.63	26	113.0650	30
AUG 23	6.864.7924	7	0.004881	13	454.77	9	80.1580	14	165.8943	9	356.30	46	32.9843	18

Table 2(concl'd.) Orbital parameters of Ariel 3

REFERENCES

<u>No.</u>	<u>Author</u>	<u>Title, etc.</u>
1	-	A discussion on scientific results obtained by the Ariel III satellite. Proc. Roy. Soc. A, <u>311</u> , 477-604 (1969)
2	R.H. Gooding	The orbit of Ariel 2 (1964-15A) - the first twelve months. R.A.E. Technical Report 65274 (1965)
3	R.H. Gooding R.J. Tayler	A PROP3 users' manual. R.A.E. Technical Report 68299 (1968)
4	R.H. Gooding	Orbit determination from Minitrack observations. R.A.E. Technical Report 66360 (1966) Phil. Trans. Roy. Soc. A, <u>262</u> , 79-88 (1967)
5	E.R. Watkins, Jr.	Preprocessing of Minitrack data. NASA Technical Note D-5042 (1969)
6	J. Hewitt	The 24 in. Schmidt satellite cameras, and their use in geodetic and geophysical investigations. Phil. Trans. Roy. Soc., A, <u>262</u> 26-31 (1967)
7	R.H. Merson	The dynamic model for PROP, a computer program for the refinement of the orbital parameters of an earth satellite. R.A.E. Technical Report 66255 (1966)
8	R.H. Gooding	A PREP users' manual. R.A.E. Technical Report 69104 (1969)
9	G.E. Cook	Variations in exospheric density during 1967-8, as revealed by ECHO 2. R.A.E. Technical Report 69127 (1969)
10	D.G. King-Hele Diana W. Scott	Further determinations of upper-atmospheric rotational speed from analysis of satellite orbits. R.A.E. Technical Report 67179 (1967) Planet. Space Sci., <u>15</u> , 1913 (1967)

REFERENCES (Contd)

<u>No.</u>	<u>Author</u>	<u>Title, etc.</u>
11	Myrna M. Lewis	Perturbations of satellite orbits by the gravitational attraction of a third body. R.A.E. Technical Report 65118 (1965)
12	Y. Kozai	Effect of precession and nutation on the orbital elements of a close earth satellite. Astronom. J., <u>65</u> , 621-623 (1960)
13	D.G. King-Hele G.E. Cook Diana W. Scott	Evaluation of odd zonal harmonics in the geo-potential, of degree less than 33, from the analysis of 22 satellite orbits. R.A.E. Technical Report 68202 (1968) Planet. Space Sci., <u>17</u> , 629-664 (1969)

004 902144

Fig.1

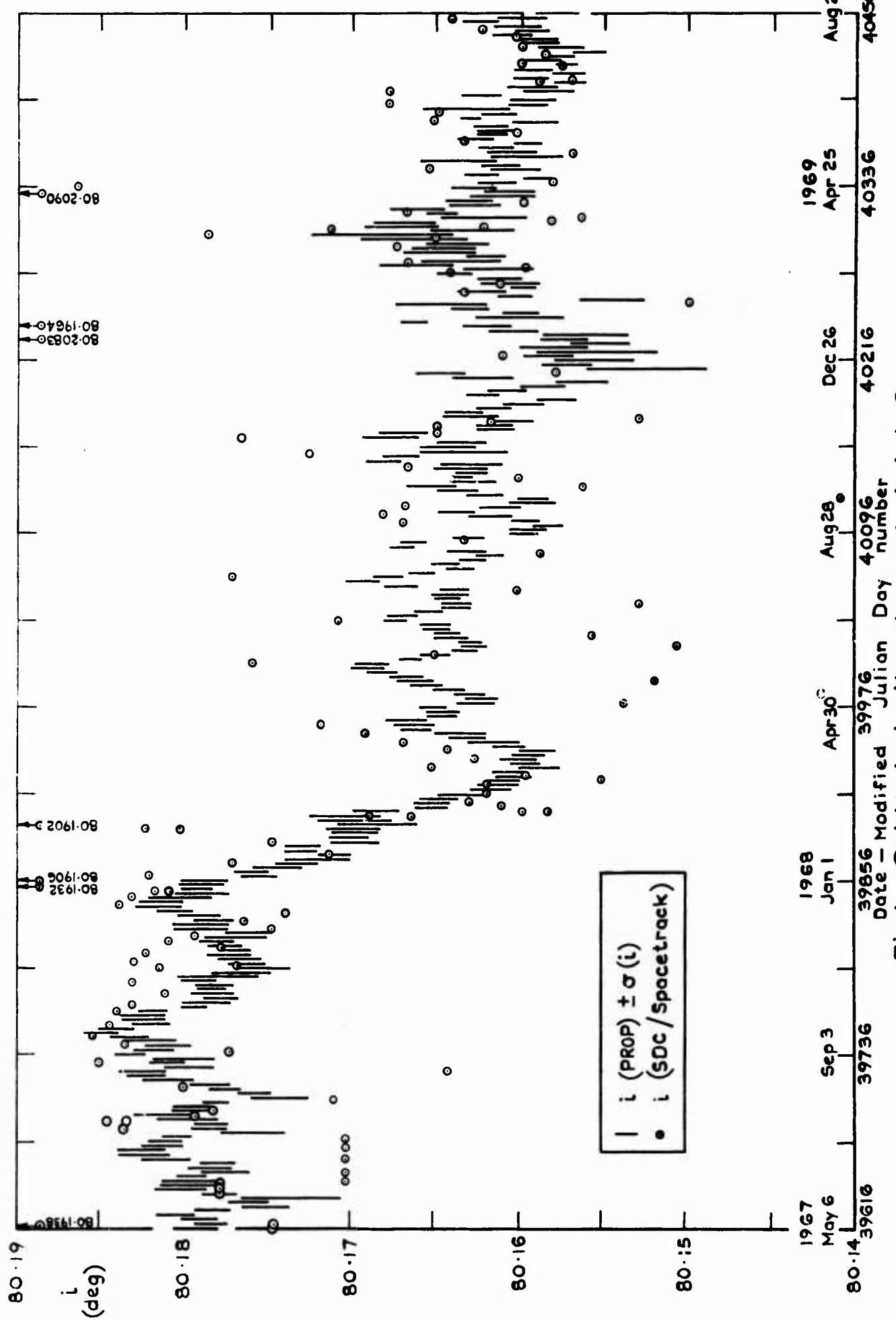


Fig. 2

004 902145

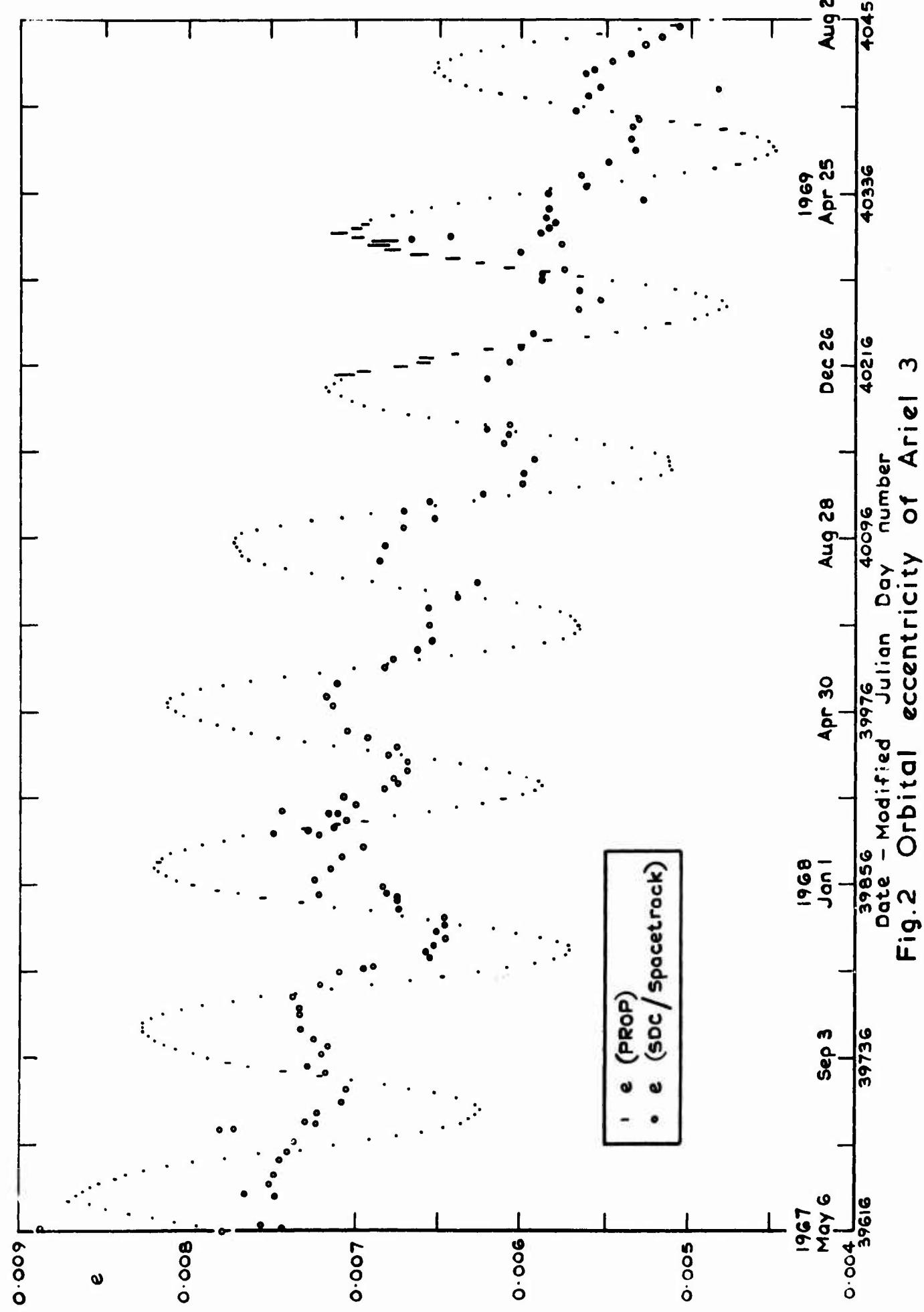


Fig. 2 Orbital eccentricity of Ariel 3

1967 May 6 39616
1968 Jan 1 39736
1969 Aug 23 40096
1969 Apr 25 40216
40456

Date - Modified Julian Day number

004 902148

Fig. 3

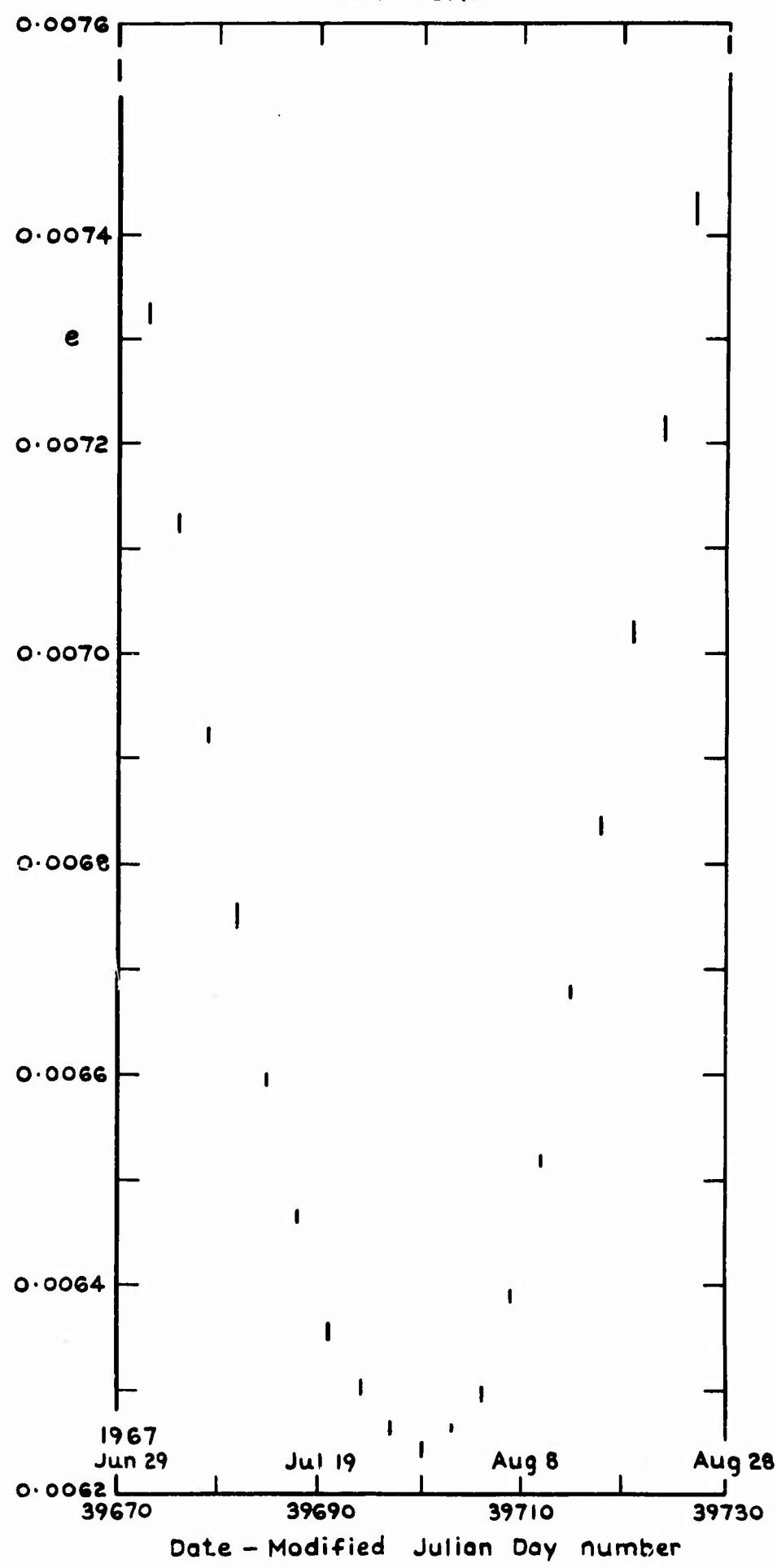


Fig. 3 Eccentricity of Ariel 3 during 60 days, on expanded scale

40 cys.

AD-705024

ROYAL AIRCRAFT ESTABLISHMENT

Technical Report 69275

December 1969

THE ORBIT OF ARIEL 3 (1967-42A)

by

R. H. Gooding

ADDENDUM

The striking change of i (orbital inclination), from 80.18° in December 1967 to about 80.163° in March 1968, was remarked upon in section 6, but left unexplained. It is believed that the explanation is now known. During this period Ariel 3 was passing through a resonance associated¹⁴ with the earth's tesseral harmonics of order 15. In fact the mean motion of the satellite was exactly commensurate with the earth's rotation rate just before 0 hours on MJD 39889 (3 February 1968), and the variable $M + \omega + 15(\Omega - n_o t)$ of Ref. 14 varied by no more than 120° from its resonance value, during the three-month period. The author, in section 6, dismissed resonance, thinking of luni-solar resonance rather than tesseral-harmonic resonance, and discussed alternative explanations which must now be considered irrelevant.

ERRATA

Page 7, line 3: for the first ' M_1 ' read ' M_1^2 '.

Page 14, line 13: for the second ' i ' read ' M_1 '.

705024

REFERENCES (Contd)

<u>No.</u>	<u>Author</u>	<u>Title, etc.</u>
14	R. R. Allan	Resonance effects due to the longitude dependence of the gravitational field of a rotating primary. R.A.E. Technical Report 66279 (1966) Planet. Space Sci., <u>15</u> , 53-76 (1967)

RECEIVED

JUL 27 1970

INPUT BY
STEARNSHOUSE